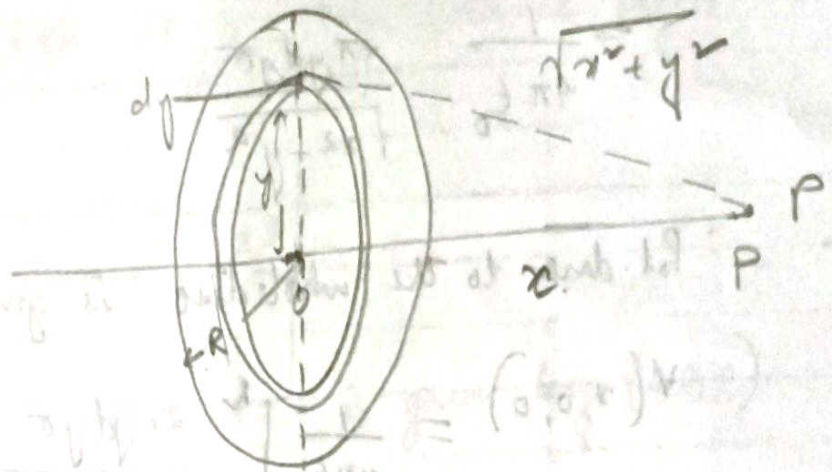


Q. 1. Find the electric potential at a point on the axis of a uniformly charged disc of surface charge density σ and radius R .

8

Ans:



Let us consider a disc lying in the $y-z$ plane. 'P' be a pt. lying on the x -axis at a distance x from the centre. x -axis has been taken as the axis of symmetry of the disc. The disc may be supposed to be formed by a large number of thin concentric ring shaped elements. Consider one such ring of radius r and thickness dy .

The area of the ring element = $2\pi r dy$

The charge on this element, $dq = 2\pi r dy \sigma$; [$\sigma \times dA$]

[σ is the surface charge density]

dist. of all points lying on the ring from P is $\sqrt{r^2 + y^2}$

∴ Potential at P due to the ring is

$$dV = \frac{1}{4\pi\epsilon_0} \frac{dq}{\sqrt{r^2 + y^2}}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{2\pi y dy \sigma}{\sqrt{r^2 + y^2}}$$

∴ Pot. due to the whole disc is given by

$$V(x, 0, 0) = \frac{1}{4\pi\epsilon_0} \int_0^R \frac{2\pi y dy \sigma}{\sqrt{r^2 + y^2}}$$

$$= \frac{2\pi\sigma}{4\pi\epsilon_0} \int_0^R \frac{y dy}{\sqrt{r^2 + y^2}}$$

$$= \frac{\sigma}{2\epsilon_0} \left[(r^2 + y^2)^{1/2} \right]_0^R$$

$$= \frac{\sigma}{2\epsilon_0} \left[(R^2 + r^2)^{1/2} - r \right]$$

This gives the potential at a point on the axis of a uniformly charged disc of charge density σ

Q. Show graphically how the electric potential varies with distance of the point from the centre.

Ans: We have, the expression for electric potential at a point on the axis at a certain distance from the centre as

$$V = \frac{\sigma}{2\epsilon_0} \left[(R^2 + z^2)^{1/2} - z \right]$$

At the centre of the disc, $z=0$ is given by $(0,0,0)$.

$$= \frac{\sigma}{2\epsilon_0} R \left[\frac{R}{R} + \frac{z}{R} \right] = \frac{\sigma R}{2\epsilon_0} \left[1 + \frac{z}{R} \right]$$

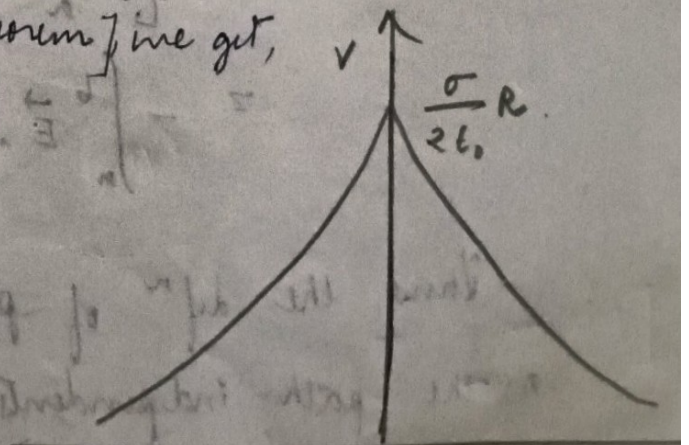
The potential goes on decreasing when we move away from the centre of the disc and at a point sufficiently far away it is given by

$$\frac{\sigma z}{2\epsilon_0} \left[\left(1 + \frac{R^2}{z^2} \right)^{1/2} - 1 \right] = \frac{\sigma z}{2\epsilon_0} \left[1 + \frac{1}{2} \frac{R^2}{z^2} - 1 \right]$$

[By applying Binomial theorem] we get,

$$V = \frac{\sigma z}{4\epsilon_0} \frac{R^2}{z^2} = \frac{1}{4\epsilon_0} \frac{\sigma R^2}{z}$$

The graph betⁿ z and V is shown as



✓ Establish that, for an electric field due to a static charge, $\boxed{\vec{E} = -\nabla V}$

Ans:

Let us define the electric potential at a local P as $V(P) = -\int_0^P \vec{E} \cdot d\vec{l}$ where 0 is a certain origin

[by convention, origin is taken at infinity]

Then the pot. diff. betⁿ two pts a & b can be replaced as

$$V(b) - V(a) = -\int_0^b \vec{E} \cdot d\vec{l} + \int_0^a \vec{E} \cdot d\vec{l} \quad (2)$$

$$= -\int_0^b \vec{E} \cdot d\vec{l} - \int_a^0 \vec{E} \cdot d\vec{l}$$

$$= -\left[\int_a^0 \vec{E} \cdot d\vec{l} + \int_0^b \vec{E} \cdot d\vec{l} \right]$$

$$= -\int_a^b \vec{E} \cdot d\vec{l}$$

→ (2)

Thus, the defⁿ of potential difference is based on the path independent of the integration in

The RHS of eqn (2). In other words, it is due to the cond

$$\text{that } \boxed{\vec{\nabla} \times \vec{E} = 0}$$

Now, from the divergence theorem of gradient, we have,

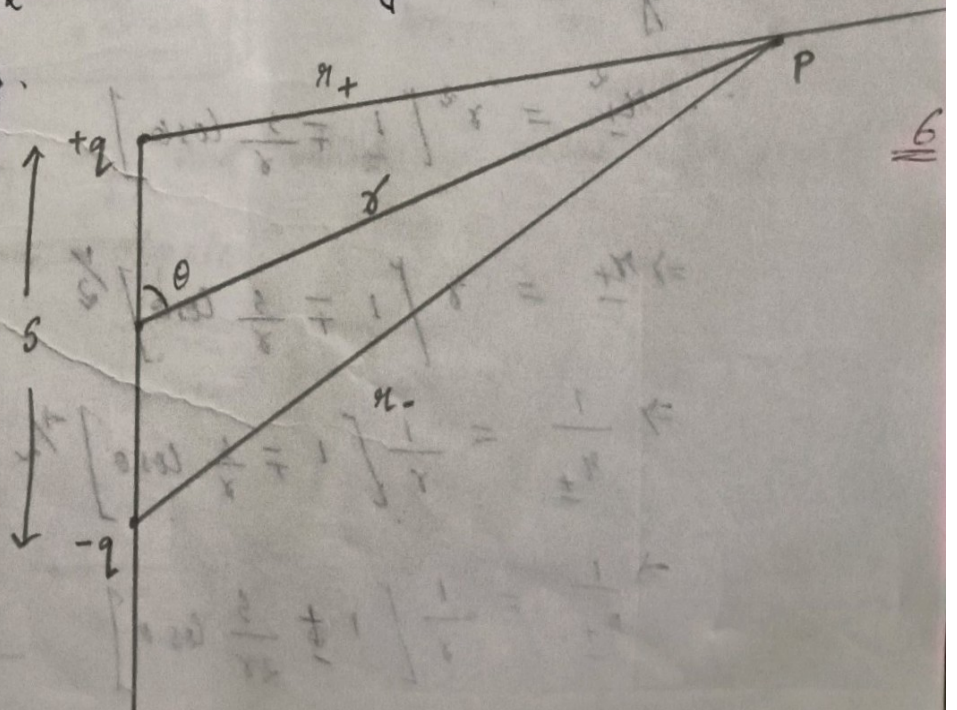
$$V(b) - V(a) = \int_a^b (\vec{\nabla} V) \cdot d\vec{l} \quad \text{--- (3)}$$

Equating (2) & (3), we get

$$\boxed{\vec{E} = -\vec{\nabla} V}$$

2. Show that the potential due to an electric dipole is $V(P) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \vec{r}}{r^2}$ where the symbols denote their usual meaning.

Ans:



In the figure, P is an arbitrary point near an electric dipole. The relevant parameters have been explained in the figure.

The electric potential at P on account of the dipole is

$$V(P) = \frac{1}{4\pi\epsilon_0} \left[\frac{q}{r_+} - \frac{q}{r_-} \right] \longrightarrow (1)$$

With the figures,

$$r_+^2 = r^2 + \left(\frac{s}{2}\right)^2 - 2r \frac{s}{2} \cos \theta$$

$$= r^2 + \frac{s^2}{4} - rs \cos \theta$$

$$= r^2 \left[1 + \frac{s^2}{4r^2} - \frac{s}{r} \cos \theta \right] \longrightarrow (2)$$

In the context of a dipole $r \gg s$, so the term $\frac{s^2}{4r^2}$ can be neglected.

$$\therefore r_+^2 = r^2 \left[1 - \frac{s}{r} \cos \theta \right]$$

$$\Rightarrow r_+ = r \left[1 - \frac{s}{r} \cos \theta \right]^{\frac{1}{2}}$$

$$\Rightarrow \frac{1}{r_+} = \frac{1}{r} \left[1 - \frac{s}{r} \cos \theta \right]^{-\frac{1}{2}}$$

$$\Rightarrow \frac{1}{r_+} = \frac{1}{r} \left[1 + \frac{s}{2r} \cos \theta \right] \longrightarrow (3)$$

[making a binomial expansion, retaining only the first 2 terms]

$$\therefore \frac{s}{r} \ll 1$$

$$\therefore \frac{1}{r_+} - \frac{1}{r_-} = \frac{1}{r} \left[1 + \frac{s}{r} \cos \theta \right] - \frac{1}{r} \left[1 - \frac{s}{r} \cos \theta \right]$$

$$\frac{1}{r_+} - \frac{1}{r_-} = \frac{2s \cos \theta}{r^2} \longrightarrow (4)$$

Substituting (4) in eqⁿ (1), we get

$$V(r) = \frac{q}{4\pi\epsilon_0} \frac{2s \cos \theta}{r^2}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{p \cos \theta}{r^2} \quad [qs = p = \text{dipole moment}]$$

$$V(r) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \hat{r}}{r^2}$$

$$= \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \vec{r}}{r^3}$$

which is the reqd exprⁿ and which is valid for points far away from a dipole.