

Deduce Kepler's third law from Newton's law of gravitation.

To a first approximation, the orbit of the planet can be considered to be circular.

Let us consider a planet of mass m moving around the sun of mass M in a circular orbit of radius r with a uniform angular velocity ω . Then, for equilibrium,

$$G \frac{Mm}{r^2} = mr\omega^2$$

$$\Rightarrow GM = r^3\omega^2$$

$$\Rightarrow GM = r^3 \cdot \frac{u\pi^2}{T^2}$$

$$\Rightarrow T^2 = \frac{u\pi^2}{GM} \cdot r^3$$

As, $\frac{u\pi^2}{GM} = \text{constant}$,

So, $T^2 \propto r^3$.

This is Kepler's third law.

Deduce Newton's law of gravitation from Kepler's law.

Let a planet of mass m be revolving around the sun of mass M in a nearly circular orbit of radius r with a uniform angular velocity ω . If T is the period of revolution of the planet, then

$$\omega = \frac{2\pi}{T}$$

Now, the centripetal force acting on the planet is,

$$F = mr\omega^2 = mr \frac{4\pi^2}{T^2} \quad \dots(1)$$

According to Kepler's 3rd law

$$T^2 \propto r^3$$

$$T = kr^3$$

Where K is constant of proportionality.

$$\therefore F = mr \frac{4\pi^2}{kr^3} = \frac{4\pi^2 m}{kr^2} \quad \dots (2)$$

$$\text{As } \frac{4\pi^2 m}{K} = \text{Constant, so, } F \propto \frac{1}{r^2}$$

This centripetal force is provided by the gravitational attractive exerted by the planet on sun. As gravitational force is a mutual force and as it depends on the mass of the planet, so it must also depend on the mass of sun M .

$$\frac{4\pi^2}{K} \propto M$$

$$\text{i.e., } \frac{4\pi^2}{K} = GM$$

Where G is the constant of proportionality. Putting this value in eqn. (2), we get,

$$F = G \frac{Mm}{r^2}$$

This is Newton's law of gravitation.

Deduce Kepler's second law.

3. If $\frac{dA}{dt}$ is the area swept out by the radius vector of a particle per unit time i.e., if $\frac{dA}{dt}$ is the areal velocity of a particle, then the angular momentum of the particle of mass m is given by,

$$L = 2m \frac{dA}{dt} \quad \dots (1)$$

Let a planet of mass m be revolving around the sun in a circular orbit of radius r . The force of attraction \vec{F} on the planet due to the sun acts along the line joining the centre of the sun and the planet and is directed towards the sun. So, the angle between \vec{F} and \vec{r} is 180° .

$$\text{Torque on the planet} = \tau = \vec{r} \times \vec{F} = rF \sin 90^\circ$$

$$\tau = rF \sin 180^\circ = 0$$

Now,

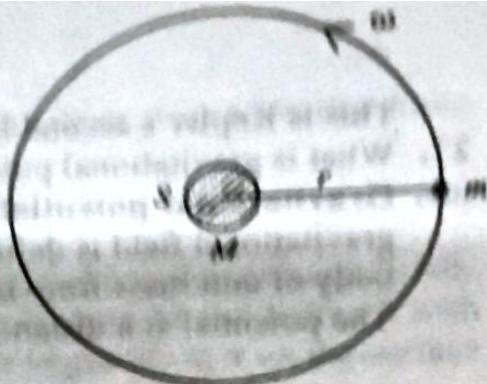
$$\tau = \frac{dL}{dt}$$

$$\therefore \frac{dL}{dt} = 0$$

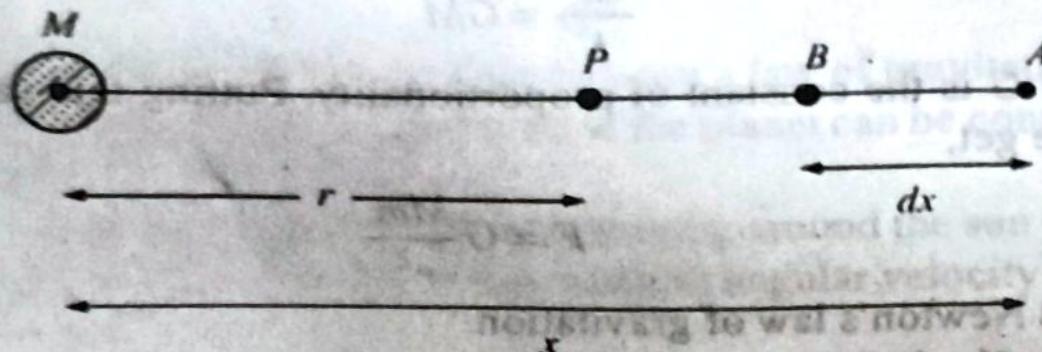
$$L = \text{constant.}$$

Hence from eqn. (1),

$$2m \frac{dA}{dt} = \text{Constant}$$



- Find the gravitational potential in the gravitational field of earth.
 Let M be the mass and R be the radius of earth. The mass of the earth can be supposed to be concentrated at its centre.
 Let us find the potential at P at a distance r from the centre of the earth ($r > R$).



The force on a unit mass at A at a distance x from the centre of the earth is,

$$F = G \frac{M}{x^2}$$

The work done in moving the unit mass through a small distance dx from A to B is given by,

$$dW = F dx = G \frac{M}{x^2} dx.$$

Hence the total work done in bringing the unit mass from infinity to P , i.e., the potential at P is given by,

$$V = W = \int_{\infty}^r \frac{GM}{x^2} dx$$

$$= \int_{\infty}^r GMx^{-2} dx$$

$$= -GM \left[\frac{1}{x} \right]_{\infty}^r$$

$$= -GM \left[\frac{1}{r} - \frac{1}{\infty} \right]$$

$$= -\frac{GM}{r}$$

$$V = -\frac{GM}{r}$$

The negative sign means that gravitational force is an attractive force.
Potential on the surface of the earth is

$$V_s = -G \frac{M}{R}$$

Find the value of acceleration due to gravity above the surface of earth.
Let us consider the earth to be a sphere of mass M and radius R with centre at O . If g_0 is the acceleration due to gravity at A on the surface of the earth, then,

$$g_0 = G \frac{M}{R^2} \quad \dots(1)$$

If g_h is the acceleration due to gravity at the point B at a height h above the surface of the earth, then,

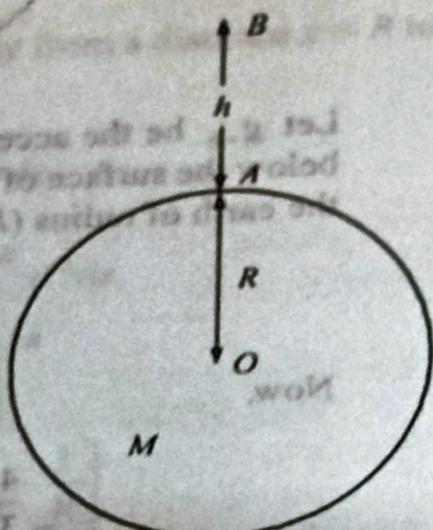
$$g_h = G \frac{M}{(R+h)^2} \quad \dots(2)$$

From (1) and (2),

$$\frac{g_h}{g_0} = \frac{R^2}{(R+h)^2}$$

$$\Rightarrow \frac{g_h}{g_0} = \frac{R^2}{R^2 \left(1 + \frac{h}{R}\right)^2}$$

$$\Rightarrow g_h = \frac{g_0}{\left(1 + \frac{h}{R}\right)^2} \quad \dots(3)$$



$$\therefore g_h < g_0$$

So, acceleration due to gravity decreases with altitude.

$$\text{From (1), } g_h = g_0 \left(1 + \frac{h}{R}\right)^{-2}$$

If $h \ll R$, then,

$$g_h = g_0 \left(1 - \frac{h}{R}\right).$$

Find the value of acceleration due to gravity below the surface of earth.
Let us consider the earth to be a homogeneous sphere of mass M and radius R with centre at O . If g is the acceleration due to gravity at the point A on the surface of the earth, then, $g = G \frac{M}{R^2}$.

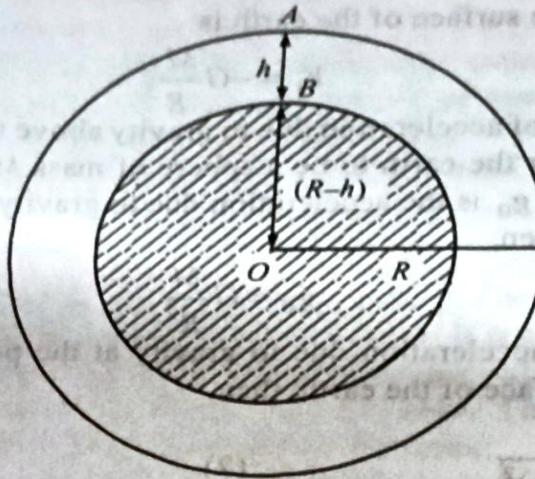
If ρ is the density of earth, then

$$M = \frac{4}{3} \pi R^3 \rho$$

$$g = \frac{G \frac{4}{3} \pi R^3 \rho}{R^2}$$

i.e.,

$$g = \frac{4}{3}\pi G R \rho$$



Let g_{-h} be the acceleration due to gravity at the point B at a depth below the surface of earth. The body will be attracted by the portion of the earth of radius $(R - h)$ and mass M' .

$$g_{-h} = G \frac{M'}{(R-h)^2}$$

$$M' = \frac{4}{3}\pi(R-h)^3\rho$$

$$\therefore g_{-h} = G \frac{\frac{4}{3}\pi(R-h)^3\rho}{R-h} = \frac{4}{3}\pi G(R-h)\rho \quad \dots(2)$$

From (1) and (2)

$$\frac{g_{-h}}{g} = \frac{\frac{4}{3}\pi G(R-h)\rho}{\frac{4}{3}\pi G R \rho} = \frac{R-h}{R} = 1 - \frac{h}{R}$$

$$\therefore g_{-h} = g \left(1 - \frac{h}{R}\right)$$

i.e.,

$$g_{-h} < g$$

Thus, acceleration due to gravity decreases with depth.

If $h = R$, i.e., at the centre of the earth,

$$g_{-h} = g(1-1) = 0$$

So, the acceleration due to gravity at the centre of the earth is zero. Hence a body becomes weightless at the centre of the earth.

State the laws of falling bodies.

The followings are the laws of falling bodies.

- (i) All bodies, falling freely from rest in vacuum, fall with equal rapidity.
- (ii) The velocity acquired by a body, falling freely from rest in vacuum, is directly proportional to the time of fall,
i.e., $v \propto t$.

(iii) The distance through which a body falls freely from rest in vacuum is directly proportional to the square of the time of fall.

$$i.e., h \propto t^2$$

Find the work done in taking a body from the surface of the earth to a height equal to the radius of the earth.

The force of gravity on a body of mass m at a distance x from the centre of the earth of mass M and radius R is given by,

$$F = G \frac{Mm}{x^2}$$

where, G = universal gravitational constant.

The work done in taking the body through a small distance dx is,

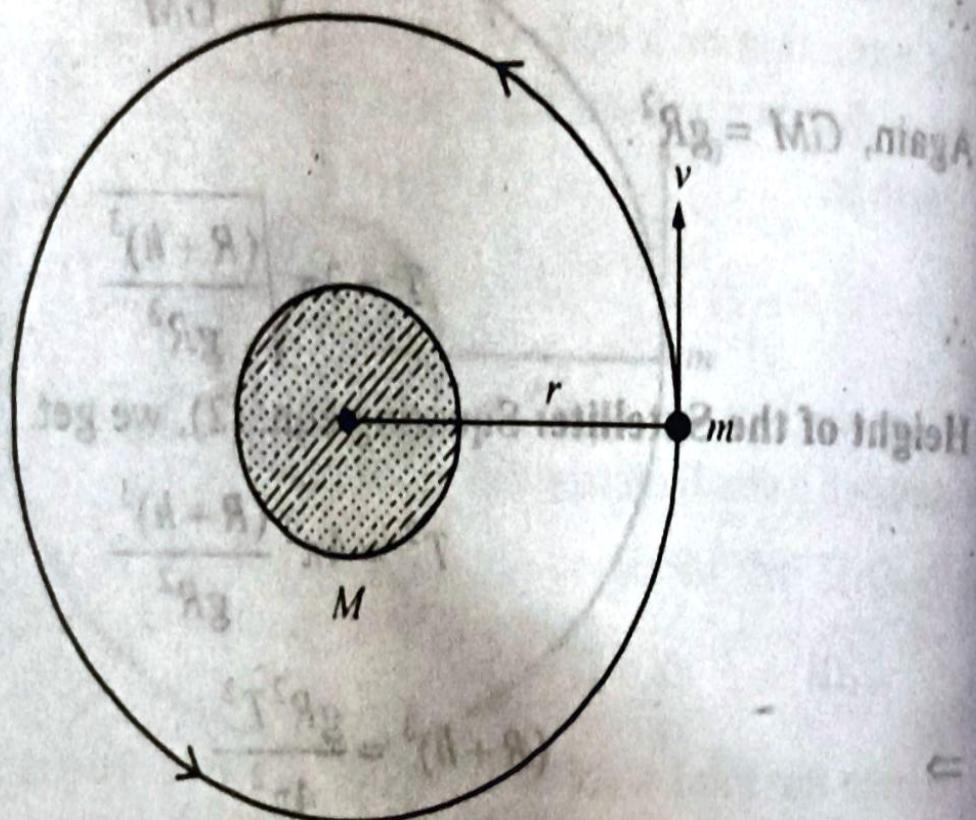
$$dW = -F dx = -G \frac{Mm}{x^2} dx$$

Hence, the total work done in taking the body from a distance $x = R$ to a distance $x = 2R$ is,

$$\begin{aligned} W &= - \int_R^{2R} G \frac{Mm}{x^2} dx \\ &= -G Mm \int_R^{2R} x^{-2} dx \\ &= G Mm \left[\frac{1}{x} \right]_R^{2R} \\ &= G Mm \left[\frac{1}{2R} - \frac{1}{R} \right] \end{aligned}$$

$$= -G \frac{Mm}{2R}$$

5. Find the energy of a satellite revolving around the earth.
- s. Let a satellite of mass m be revolving around the earth of mass M with an orbital velocity v in a circular orbit of radius r . At any instant, the satellite possesses two types of energy.
- (i) **Kinetic energy:** It is the energy possessed by the satellite due to motion. Kinetic energy is given by,



$$T = \frac{1}{2}mv^2$$

Again

$$\frac{mv^2}{r} = G \frac{Mm}{r}$$

i.e.,

$$mv^2 = G \frac{Mm}{r}$$

$$\therefore T = \frac{1}{2} \frac{G Mm}{r} \quad \dots(1)$$

(ii) **Potential energy:** It is the energy possessed by the satellite due to its position in the orbit.

Potential energy

$$= U = -\frac{G Mm}{r} \quad \dots(2)$$

∴ Total energy

$$\begin{aligned} &= E = T + U \\ &= \frac{1}{2} \frac{G Mm}{r} - \frac{G Mm}{r} \\ &= -\frac{1}{2} \frac{G Mm}{r} \end{aligned}$$

The negative sign means that the satellite is bound to the planet.